

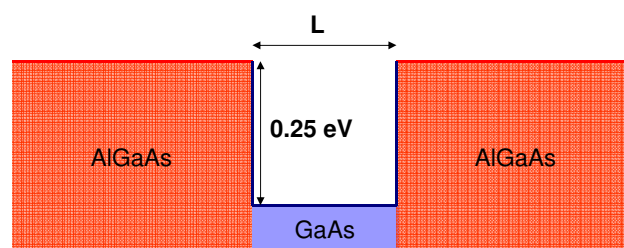
Simulating the Energy Spectrum of Quantum Dots

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PROBLEM 1. Calculate the electron energy spectrum of a 1D GaAs/AlGaAs QD as a function of the size.

Hint: consider GaAs effective mass all over the structure.



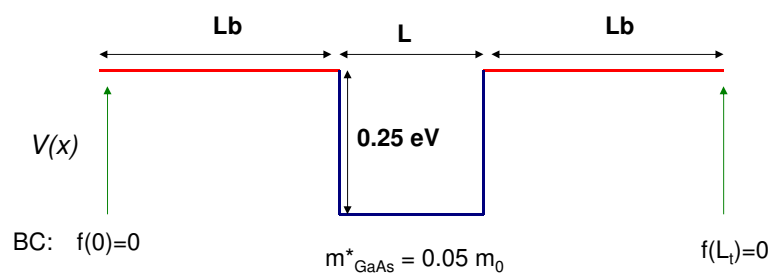
$$m_{\text{GaAs}}^* = 0.05 m_0$$

The single-band effective mass equation:

$$\left[-\frac{\hbar^2}{2m^*} \frac{d^2}{dx^2} + V(x)\right] f(x) = E f(x)$$

Let us use atomic units ($\hbar=m_0=e=1$)

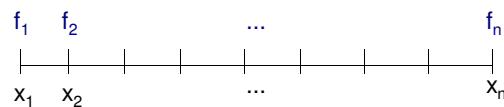
$$\left[-\frac{1}{2(m^*/m_0)} \frac{d^2}{dx^2} + V(x)\right] f(x) = E f(x)$$



Numerical integration of the differential equation: *finite differences*

$$\left[-\frac{1}{2m^*} \frac{d^2}{dx^2} + V(x)\right] f(x) = E f(x)$$

Discretization grid

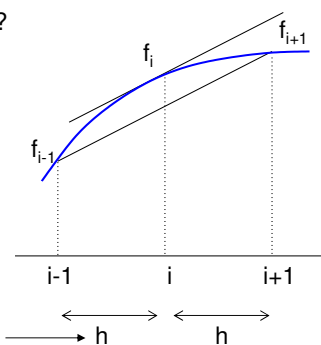


How do we approximate the derivatives at each point?

$$f'(x_i) = f'_i = \frac{f_{i+1} - f_{i-1}}{2h}$$

$$f''(x_i) = f''_i = \frac{f_{i+1} - f'_i}{2h} =$$

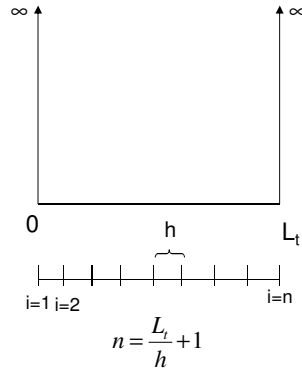
$$= \dots = \frac{f_{i+1} - 2f_i + f_{i-1}}{h^2}$$



Step of the grid

FINITE DIFFERENCES METHOD

$$\left[-\frac{1}{2m^*} \frac{d^2}{dx^2} + V(x)\right] f(x) = E f(x) \quad + \quad BCs: \begin{cases} f(0)=0 \\ f(L_t)=0 \end{cases}$$



1. Define discretization grid

2. Discretize the equation:

$$-\frac{1}{2m^*} f_i'' + V_i f_i = E f_i$$

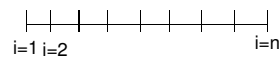
$$-\frac{1}{2m^* h^2} [f_{i+1} - 2f_i + f_{i-1}] + V_i f_i = E f_i$$

3. Group coefficients of fwd/center/bwd points

$$\left(-\frac{1}{2m^* h^2}\right) f_{i-1} + \left(\frac{1}{m^* h^2} + V_i\right) f_i + \left(-\frac{1}{2m^* h^2}\right) f_{i+1} = E f_i$$

$$\beta_{i-1} f_{i-1} + a_i f_i + b_{i+1} f_{i+1} = E f_i$$

$$\beta_{i-1} f_{i-1} + a_i f_i + b_{i+1} f_{i+1} = E f_i$$



Trivial eqs: $f_1 = 0$, $f_n = 0$.

Extreme eqs:

$$\begin{aligned} i=2 &\rightarrow \beta_1 \overset{0}{f_1} + a_2 f_2 + b_3 f_3 = E f_2 \\ \dots &\dots \dots \dots \dots \\ i=n-1 &\rightarrow \beta_{n-2} f_{n-2} + a_{n-1} f_{n-1} + b_n \underset{0}{f_n} = E f_{n-1} \end{aligned}$$

Matriz (n-2) x (n-2) - sparse

We now have a standard diagonalization problem (dim n-2):

$$\begin{bmatrix} a_2 & b_3 & & & \\ \beta_2 & a_3 & b_4 & & \\ & \ddots & \ddots & \ddots & \\ & & \ddots & a_{n-2} & b_{n-1} \\ & & & \beta_{n-2} & a_{n-1} \end{bmatrix} \cdot \begin{bmatrix} f_2 \\ f_3 \\ \vdots \\ f_{n-2} \\ f_{n-1} \end{bmatrix} = E \begin{bmatrix} f_2 \\ f_3 \\ \vdots \\ f_{n-2} \\ f_{n-1} \end{bmatrix}$$

```

% INPUT DATA

L_nm=40.0;      % box length (nm)
Lb_nm=20.0;     % length of the barrier (nm)
m=0.05;         % electron effective mass (times m0)
Vin_eV=0.00;    % potential energy in the box (eV)
Vout_eV=0.1;     % potential energy in the barrier (eV)
n=321;          % grid nodes
neig=4;         % number of states to be calculated (starting from ground state)

% 1) Convert into atomic units
kdist=0.0529177; % kdist nm = 1 bohr radius
kener=27.21138;  % kener eV = 1 hartree

L=L_nm/kdist;
Lb=Lb_nm/kdist;
Vin=Vin_eV/kener;
Vout=Vout_eV/kener;

% 2) Define grid

Lt= 2*Lb + L;    % total length
h=Lt/(n-1)       % grid step

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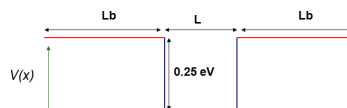
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% 3) Calculate potential V(x)

for i=1:n
    x(i)= (i-1)*h-Lt/2; % position at i-th grid point
    if (abs(x(i))<=L/2)
        V(i)=Vin; % inside the box
    else
        V(i)=Vout; % in the barrier
    end;
end;

% 4) Define coefficients of f(i), f(i+/-1)
a=zeros(n,1);
for i=1:n
    a(i)=1/(m*h^2)+V(i); % coef of f(i) -depends on i
end;
b=-1/(2*m*h^2);        % coef. of f+ and f-

```



$$\begin{bmatrix}
 a_2 & b & & & \\
 b & a_3 & b & & \\
 & \ddots & \ddots & \ddots & \\
 & & b & a_{n-2} & b \\
 & & & b & a_{n-1}
 \end{bmatrix}$$

[di	dc	ds]
2	2	X
3	3	X
X	4	4
X	5	5

```
mat=spdiags([di dc ds],[-1, 0, 1],4,4)
```

2	3	0	0
2	3	4	0
0	3	4	5
0	0	4	5

$$\beta_{i-1} f_{i-1} + a_i f_i + b_{i+1} f_{i+1} = E f_i$$

$$\begin{bmatrix} a_2 & b_3 & & & \\ \beta_2 & a_3 & b_4 & & \\ & \ddots & \ddots & \ddots & \\ & & \ddots & a_{n-2} & b_{n-1} \\ & & & \beta_{n-2} & a_{n-1} \end{bmatrix}$$

β_2	a_2	0
β_3	a_3	b_3
...	...	b_4
...	...	...
...	...	...
β_{n-2}	a_{n-2}	b_{n-2}
0	$a_{(n-1)}$	b_{n-1}

```
% 5) Build the Hamiltonian matrix

% define the three diagonals (vectors)

dc=a(2:n-1); % main diagonal f(i)
di=[ones(n-3,1)*b;0]; % or just di=ones(n-2,1)*b; subdiagonal f(i-1)
ds=[0;ones(n-3,1)*b]; % or just ds=ones(n-2,1)*b; superdiagonal f(i+1)

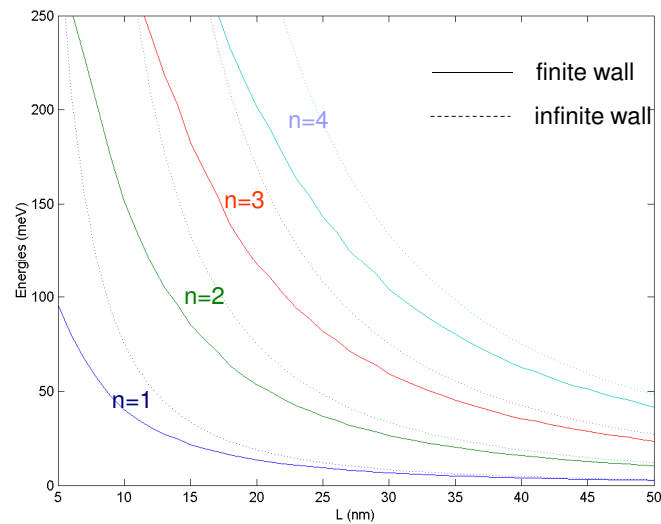
% introduce them into a sparse, tridiagonal matrix
matrix=spdiags([di,dc,ds],[-1,0,1],n-2,n-2);

% 6) Diagonalize the matrix
[P,D]=eigs(matrix,neig,'sa');
eners=diag(D)*kener; % grep energy and convert to eV

% plotting potential vs. x (nm)
figure;plot(x*kdist,V*kener);
% plotting the neig low-lying eigenfunctions vs. x (nm)
figure;plot(x(2:n-1)*kdist,P(:,1:neig));

% write neig low-lying energy meV
'Lowest energies meV ', eners
```

The result should look like this:



Exercices:

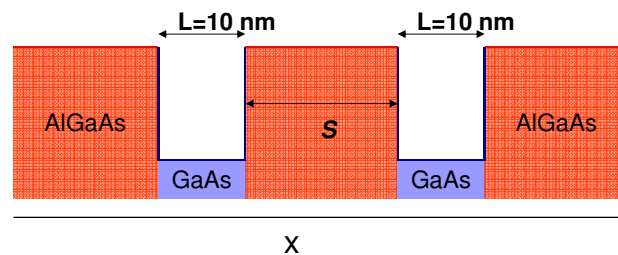
- a) Compare the converged energies with those of the particle-in-the-box with infinite walls for the $n=1,2,3$ states.

$$E_n = \frac{\hbar^2 \pi^2}{2mL^2} n^2$$

- b) Plot the 3 lowest eigenstates for $L=15$ nm, $L_b=10$ nm. What is different from the infinite wall eigenstates?

HOME WORK:

Calculate the electron energy spectrum of two coupled QDs as a function of their separation S . Plot the two lowest states for $S=1$ nm and $S=10$ nm.



The result should look like this:

